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# Efficient utilization of chemical reaction for energy Activation on MHD flow of Prandtl Nano liquid

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#### Abstract:

An examination is refined to investigate the result of energy initiation on MHD stream of Prandtl Nano liquid with compound response. The overseeing PDEs are changed into ODEs utilizing suitable change, which are mathematically addressed by means of RK scheme alongside shooting technique. Graphical outcomes clarify the impact of velocity, temperature and nanoparticles fixation. Also, drag coefficient, Nusselt and Sherwood number are thought and offered by means of table. Results show that fluid velocity upgrades by expanding Prandtl boundary just as flexible boundary, while invert exercises are noted for temperature and nanoparticles fixation profiles.

#### **Keywords:**

MHD, Rk, Nano, ODE, PDE, Activation, chemical.

#### Introduction:

These days, the utilization of nanotechnology is exceptionally wide in enterprises. Among other numerous uses of nanoscience, nanofluid is the developing field in heat move, hardware, aviation, cars, sun oriented energy frameworks, drug measures and so on A suspension of nanoparticles is known as nanofluid and the element of nanofluid is under 100nm in the base liquids (oil, glycol, water and so on) To create liquids having suspension solidness and higher convection, Choi [1] presented the idea of nanofluids. He uncovered that with a little amassing of nanomaterials the warm conduction of base liquids could roughly be multiplied [2]. Bhatti et al. [3] offered the stagnation point magnetohydrodynamic (MHD) stream past a permeable sheet utilizing hearty numerical method. The mathematical examination of nanofluid with consistent magnetic field containing nanoparticles of copper is concentrated by Sheikholeslami and Ganji [4]. The entropy age of non-Newtonian nanofluid over an retentive sheet utilizing SLM was considered by

Bhatti et al. [5]. Maxwell [6] offered a numerical model which expresses that the warm conduction of nanomaterials depends on the warm conduction of circular bodies. Hamilton and Crosser [7] made further improvements on round particles. Khan and Pop [8] introduced the 2-D stream including nanofluid past extending plate. Makinde and Aziz [9] further investigated this examination utilizing convective conditions and showed that through convective warming, the warm layer is fundamentally influenced. The assurance of nanoparticles' warmth conductivity is finished by boundaries related to different nanomaterials sizes [10], surface charge of base liquids, shapes, focus and others. Sheikholeslami explored the highlights of warmth transmission in the progression of nanofluid with magnetic impact. The radiative MHD nanofluid stream with synthetic response and variable sheet in a porous medium is concentrated by Zhang . With stagnation-point and thermophoretic impact, Zaib contemplated shaky progression of nanofluid with warm convection and mass transmission. F. Ali inspected the shaky progression of Eyring-Powell nanofluid past a convectively warmed extended surface . On account of a few capacities in designing, innovative what's more, modern movements, numerous specialists have altogether considered the rheology of non-Newtonian nanofluids. Scarcely any natural standards of non-Newtonian liquids are paints, toothpaste, blood, emulsion, oil and synthetically intensified sheets and so on As indicated by their actual conduct, there exist various significant classes of non-Newtonian liquids. Among others, pseudo plastic is a huge type of non-Newtonian liquids additionally named as shear diminishing liquids. At the point when shear is applied to this kind of liquid the thickness diminishes. The notable illustration of pseudo plastic liquid is human blood. The qualities of pseudo plastic non- Newtonian liquids can be obvious by Prandtl liquid model. For Prandtl liquid, Nadeem et al. dissected the progression of blood inside the supply routes and found that the thermophoresis boundary and stream rate have reverse connection. With convective warming Soomro et al. thought the Prandtl liquid stream over an extending or contracting plate. They inferred that with the addition in Prandtl liquid boundary extending velocity increments and contracting velocity diminishes.



Fig.1 problem diagram

#### 2. Methodology of Mathematical Formulation:

$$\tau = \frac{\Gamma arc \sin \left(\frac{1}{\Omega} \left(\left(\frac{\partial \hat{u}}{\partial y}\right)^2 + \left(\frac{\partial \bar{v}}{\partial x}\right)^2\right)^{1/2}\right)}{\left(\left(\frac{\partial \hat{u}}{\partial y}\right)^2 + \left(\frac{\partial \bar{v}}{\partial x}\right)^2\right)^{1/2}} \frac{\partial \hat{u}}{\partial y}, \qquad (1)$$

$$\begin{aligned} \frac{\partial \hat{u}}{\partial x} + \frac{\partial \bar{v}}{\partial y} &= 0, \end{aligned} (2) \\ \hat{u} \frac{\partial \hat{u}}{\partial x} + \bar{v} \frac{\partial \hat{u}}{\partial y} &= v \frac{\Gamma}{\Omega} \frac{\partial^2 \hat{u}}{\partial y^2} + \\ \frac{v\Gamma}{2\Omega^3} \left(\frac{\partial \hat{u}}{\partial y}\right)^2 \frac{\partial^2 \hat{u}}{\partial y^2} - \frac{\sigma B^2}{\rho} \hat{u}, \end{aligned} (3) \\ \hat{u} \frac{\partial \bar{T}}{\partial x} + \bar{v} \frac{\partial \bar{T}}{\partial y} &= \alpha \frac{\partial^2 \bar{T}}{\partial y^2} + \\ \tau \left( D_B \frac{\partial \bar{T}}{\partial y} \cdot \frac{\partial \hat{C}}{\partial y} + \frac{D_T}{\bar{T}_{\omega}} \left( \frac{\partial \bar{T}}{\partial y} \right)^2 \right), \end{aligned} (4) \\ \hat{u} \frac{\partial \hat{C}}{\partial x} + \bar{v} \frac{\partial \hat{C}}{\partial y} &= D_B \frac{\partial^2 \hat{C}}{\partial y^2} + \frac{D_T}{\bar{T}_{\omega}} \left( \frac{\partial^2 \bar{T}}{\partial y^2} \right) - \\ K_r^2 \left( \frac{\bar{T}}{\bar{T}_{\omega}} \right)^n Exp \left( -\frac{E_a}{k\bar{T}} \right) (\hat{C} - \hat{C}_{\omega}), \end{aligned} (5) \\ \hat{u} &= \hat{u}_w = c_0 x, \ \bar{v} = 0, \quad -k \frac{\partial \bar{T}}{\partial y} = hs(x) \bar{T}, \ \hat{C} = \hat{C}_w \ \text{at } y = 0, \\ \hat{u} \to 0, \ \bar{T} \to \bar{T}_{\omega}, \ \hat{C} \to \hat{C}_{\omega} \ \text{as } y \to \infty. \end{aligned} (6) \\ \text{The expressions of transformation [18] are as follows:} \end{aligned}$$

$$\eta = \left(\frac{c_0}{\nu}\right)^{\frac{1}{2}} y, \quad \psi = (c_0 \nu)^{\frac{1}{2}} xf(\eta),$$

$$\hat{u} = c_0 xf'(\eta), \quad \tilde{\nu} = -(c_0 \nu)^{\frac{1}{2}} f(\eta),$$

$$\tilde{T} = \tilde{T}_{\omega} + \left(\tilde{T}_f - \tilde{T}_{\omega}\right)\theta(\eta),$$

$$\hat{C} = \hat{C}_{\omega} + \left(\hat{C}_w - \hat{C}_{\omega}\right)\phi(\eta).$$
(7)

$$\alpha f''' - (f')^2 + ff'' + \alpha \beta (f'')^2 f''' - Mf' = 0, \qquad (8)$$

$$\theta'' + \Pr\left(f\theta' + Nb\theta'\phi' + Nt(\theta')^{2}\right) = 0, \qquad (9)$$
  
$$\phi'' + Scf\phi' + \frac{Nt}{Nb}\theta'' -$$
  
$$\lambda Sc \left(1 + \gamma\theta\right)^{n} Exp\left(-\frac{E}{1 + \gamma\theta}\right)\phi = 0, \qquad (10)$$



#### 3. Result and Discussion:

In this article, shooting technique has been utilized to resolve the ODE conditions (8) - (10) with conditions introduced in eq. (11). In Table I, we researched the conduct of skin grating f C , Nusselt number x Nu and Sherwood number x Sh as for Prandtl liquid boundary  $\alpha$  against the flexible parameter $\beta$ . Table I infers that skin grating, neighborhood Nusselt and Sherwood number all are expanding with the expansion in  $\alpha$  and  $\beta$ . In Table II, we approve the current results with the accessible outcome introduced by Soomro . The effect of Prandtl liquid boundary  $\alpha$  upon velocity, temperature and nanoparticles fixation have been exhibited in Figs. 2-4. The impact of  $\alpha$  over velocity profile is depicted by means of Fig. 2. We can see that as  $\alpha$  builds the liquid velocity increments. By upgrading  $\alpha$  the thickness of liquid lessens, this decrease in thickness causes the increase in velocity. Figs. 3 and 4 outline the diminishing habits of temperature field and nanomaterial fixation therefore with an increment in  $\alpha$ .



**Table I:** The different values of drag coefficient, Nusselt number and Sherwood number against  $\alpha$  for different values of  $\beta$  when Nt = Bi = E = 0.1,  $M = Sc = \lambda = 1$ , Pr = 2, Nb = 1.5,  $\gamma = 0.5$ , n = 0.2.

α	β	$C_f \operatorname{Re}_x^{1/2}$	$Nu_x \operatorname{Re}_x^{-1/2}$	$Sh_x \operatorname{Re}_x^{-1/2}$
3	3	1.0640	0.0605	1.1787
	4	2.1416	0.0607	1.1803
	5	3.2039	0.0609	1.1816
	6	4.2558	0.0610	1.1828
4	3	1.6781	0.0612	1.1838
	4	3.0963	0.0614	1.1852
	5	4.4988	0.0615	1.1864
	б	5.8903	0.0616	1.1874
5	3	2.3320	0.0617	1.1877
	4	4.0885	0.0618	1.1889
	5	5.8294	0.0619	1.1899
	6	7.5591	0.0620	1.1908
6	3	3.0147	0.0620	1.1906
	4	5.1081	0.0621	1.1917
	5	7.1860	0.0622	1.1926
	6	9.2528	0.0623	1.1934

#### 4. Conclusion:

The impact of initiation energy on MHD Prandtl nanofluid with initiation energy is found mathematically over an extending a field. The accompanying key marks are drawn from this article.Fluid velocity increments for the bigger upsides of Prandtl liquid boundary  $\alpha$  and versatile boundary  $\beta$  though temperature and nanomaterial focus show diminishing way. Temperature and focus have expanding impact for the more prominent upsides of M .Concentration of nanoparticle is improved by the greater E and inverse conduct for the more prominent upsides of dimensionless substance response rate consistent  $\lambda$ . Fluid temperature increments when Biot number increments.

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